

Superconvergence of Galerkin Variational Integrators

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Lagrangian mechanics

Lagrangian $\mathcal{L} : TQ \rightarrow \mathbb{R}$ on a vector space Q . **Action** $\mathfrak{S}[q] = \int_a^b \mathcal{L}(q(t), \dot{q}(t)) dt$.

A curve q is **critical** if $\left. \frac{\partial}{\partial \alpha} \mathfrak{S}[q + \alpha \delta q] \right|_{\alpha=0} = 0$ for all $\delta q : [a, b] \rightarrow Q$ with $\delta q(a) = \delta q(b) = 0$.

A curve q is critical if the **Euler-Lagrange equation** holds: $\frac{\partial \mathcal{L}}{\partial q} - \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}} = 0$

Assume the Lagrangian is non-degenerate: $\det \left(\frac{\partial^2 \mathcal{L}}{\partial \dot{q}^2} \right) \neq 0$. Then

- ▶ the Euler-Lagrange equation is a second order ODE
- ▶ the Legendre transform $TQ \rightarrow T^*Q : (q, \dot{q}) \mapsto (q, p) = \left(q, \frac{\partial \mathcal{L}}{\partial \dot{q}} \right)$ is invertible.
- ▶ the Euler-Lagrange equation is equivalent to the Hamiltonian system

$$\dot{q} = \frac{\partial \mathcal{H}(q, p)}{\partial p}, \quad \dot{p} = -\frac{\partial \mathcal{H}(q, p)}{\partial q}$$

with Hamiltonian $\mathcal{H}(q, p) = p\dot{q} - \mathcal{L}(q, \dot{q})$.

Variational integrators

Discrete Lagrange function $L : Q \times Q \times (0, \infty) \rightarrow \mathbb{R}$.

Discrete action for a discrete curve $q = (q_0, q_1, \dots, q_N)$ with step size h :

$$\mathfrak{S}_d(q) = \sum_{i=1}^N L(q_{i-1}, q_i; h).$$

Discrete Euler-Lagrange equation

$$D_2 L(q_{i-1}, q_i; h) + D_1 L(q_i, q_{i+1}; h) = 0$$

for $i \in \{1, \dots, N-1\}$, where D_1 and D_2 denote the partial derivatives of L .

Discrete EL equation implies equality of the two formulas for the discrete momentum,

$$p_i = D_2 L(q_{i-1}, q_i; h) \quad \text{and} \quad p_i = -D_1 L(q_i, q_{i+1}; h).$$

This gives a natural implementation of the discrete Euler-Lagrange equation as a one-step method $\Phi_h : (q_i, p_i) \mapsto (q_{i+1}, p_{i+1})$ which is a symplectic integrator.

Marsden, West. [Discrete mechanics and variational integrators](#). Acta Numerica, 2001.

Variational error analysis

Assume that there exists a unique smooth minimizer of the action subject to $q(0) = q_0$ and $q(h) = q_1$. The minimal value of the action is called the **exact discrete Lagrangian**:

$$L_{\text{exact}}(q_0, q_1, h) = \min_q \int_0^h \mathcal{L}(q, \dot{q}) dt.$$

The order of a variational integrator can be determined by comparing its discrete Lagrangian to the exact discrete Lagrangian.

If

$$L(q(0), q(h), h) - L_{\text{exact}}(q(0), q(h), h) = \mathcal{O}(h^{\ell+1}),$$

then the variational integrator defined L (in its symplectic form Φ_h) is of order ℓ :

$$\Phi_h(q, p) - \varphi_h(q, p) = \mathcal{O}(h^{\ell+1}),$$

where φ_h is the flow over time h of the continuous Hamiltonian system.

Marsden, West. **Discrete mechanics and variational integrators**. Acta Numerica, 2001.

Patrick, Cuell. **Error analysis of variational integrators of unconstrained Lagrangian systems**. Numerische Mathematik, 2009.

Galerkin variational integrators

Replace the space of smooth curves by a **finite dimensional space of polynomials**

$$\mathcal{P}^s([0, h], Q) = \{q : [0, h] \rightarrow Q \text{ polynomial of degree at most } s\}$$

($\mathcal{P}^s([0, h], Q)$ can be parameterized by the values of q at $s + 1$ control points.)

As discrete Lagrangian, we want

$$L(q_0, q_1; h) \approx \min_{\substack{q \in \mathcal{P}^s([0, h], Q), \\ q(0)=q_0, q(h)=q_1}} \left(\int_0^h \mathcal{L}(q(t), \dot{q}(t)) dt \right)$$

Approximation using a **quadrature** rule with points $c_i \in [0, 1]$ and weights $b_i \in \mathbb{R}$:

$$L(q_0, q_1; h) = \min_{\substack{q \in \mathcal{P}^s([0, h], Q), \\ q(0)=q_0, q(h)=q_1}} \left(h \sum_i b_i \mathcal{L}(q(hc_i), \dot{q}(hc_i)) \right)$$

Marsden, West. **Discrete mechanics and variational integrators**. Acta Numerica, 2001.

Leok, Shingel. **General techniques for constructing variational integrators**. Frontiers of Mathematics in China, 2012

Superconvergence

Previous result*

Let L be a Galerkin discretization of a Lagrangian \mathcal{L} , based on [polynomials of degree \$s\$](#) and a [quadrature rule of degree \$u\$](#) . Assume that all discrete and continuous critical curves minimize their respective actions. Then the corresponding symplectic [integrator is of order \$\min\(s, u\)\$](#) .

Numerical evidence suggested superconvergence: order of up to $2s^\dagger$.

Superconvergence has been proved for some particular cases[‡].

Our result

The corresponding symplectic integrator is of [order \$\min\(2s, u\)\$](#) .

*Hall, Leok. [Spectral variational integrators](#). Numerische Mathematik, 2015.

†Ober-Blöbaum, Saake. [Construction and analysis of higher order Galerkin variational integrators](#) Advances in Computational Mathematics, 2015.

‡Ober-Blöbaum. [Galerkin variational integrators and modified symplectic Runge-Kutta methods](#), IMA Journal of Numerical Analysis, 2017.

Superconvergence: sketch of proof

By [variational error analysis](#), it suffices to show that for any smooth q

$$L_{\text{exact}}(q(0), q(h); h) - L(q(0), q(h); h) = \mathcal{O}\left(h^{\min(2s, u)+1}\right) \quad (*)$$

Consider three curves:

- ▶ q_{EL} : minimizer of the continuous action with $q_{\text{EL}}(0) = q(0)$ and $q_{\text{EL}}(h) = q(h)$.
- ▶ \hat{q} : polynomial that agrees with q_{EL} at control points $0 = hd_0, hd_1, \dots, hd_s = h$.
- ▶ \tilde{q} : polynomial that minimizes $h \sum_i b_i \mathcal{L}(q(hc_i), \dot{q}(hc_i))$.

Then (*) is equivalent to

$$\int_0^h \mathcal{L}(q_{\text{EL}}, \dot{q}_{\text{EL}}) dt - h \sum_i b_i \mathcal{L}(\tilde{q}(hc_i), \dot{\tilde{q}}(hc_i)) = \mathcal{O}\left(h^{\min(2s, u)+1}\right)$$

or

$$\left(\int_0^h \mathcal{L}(q_{\text{EL}}, \dot{q}_{\text{EL}}) dt - \int_0^h \mathcal{L}(\hat{q}, \dot{\hat{q}}) \right) + \left(\int_0^h \mathcal{L}(\hat{q}, \dot{\hat{q}}) - h \sum_i b_i \mathcal{L}(\tilde{q}(hc_i), \dot{\tilde{q}}(hc_i)) \right) = \mathcal{O}\left(h^{\min(2s, u)+1}\right)$$

Superconvergence: sketch of proof

Since q_{EL} and \hat{q} agree at $s + 1$ control points in $[0, h]$ there holds $q_{EL} - \hat{q} = \mathcal{O}(h^{s+1})$ and $\dot{q}_{EL} - \dot{\hat{q}} = \mathcal{O}(h^s)$. Hence

$$\begin{aligned} & \int_0^h \mathcal{L}(q_{EL}, \dot{q}_{EL}) dt - \int_0^h \mathcal{L}(\hat{q}, \dot{\hat{q}}) dt \\ &= \int_0^h \left(\frac{\partial \mathcal{L}(q_{EL}, \dot{q}_{EL})}{\partial q} (q_{EL} - \hat{q}) + \frac{\partial \mathcal{L}(q_{EL}, \dot{q}_{EL})}{\partial \dot{q}} (\dot{q}_{EL} - \dot{\hat{q}}) + \mathcal{O}(h^{2s}) \right) dt \\ &= \int_0^h \left(\left(\frac{\partial \mathcal{L}(q_{EL}, \dot{q}_{EL})}{\partial q} - \frac{d}{dt} \frac{\partial \mathcal{L}(q_{EL}, \dot{q}_{EL})}{\partial \dot{q}} \right) (q_{EL} - \hat{q}) + \mathcal{O}(h^{2s}) \right) dt \\ &\quad + \left(\frac{\partial \mathcal{L}(q_{EL}, \dot{q}_{EL})}{\partial \dot{q}} (q_{EL} - \hat{q}) \right) \Big|_0^h. \end{aligned}$$

Note that $\hat{q}(0) = q_{EL}(0)$ and $\hat{q}(h) = q_{EL}(h)$ and q_{EL} solves the EL equation, so

$$\int_0^h \mathcal{L}(q_{EL}, \dot{q}_{EL}) dt - \int_0^h \mathcal{L}(\hat{q}, \dot{\hat{q}}) dt = \int_0^h \mathcal{O}(h^{2s}) dt = \mathcal{O}(h^{2s+1}).$$

Superconvergence: sketch of proof

- ▶ Using the calculus of variations we found

$$\int_0^h \mathcal{L}(q_{\text{EL}}, \dot{q}_{\text{EL}}) dt - \int_0^h \mathcal{L}(\hat{q}, \dot{\hat{q}}) dt = \mathcal{O}(h^{2s+1}).$$

- ▶ From the assumption that critical curves are minimizers, it follows that \hat{q} and \tilde{q} are close to each other. In particular,

$$\int_0^h \mathcal{L}(\hat{q}, \dot{\hat{q}}) - h \sum_i b_i \mathcal{L}(\tilde{q}(hc_i), \dot{\tilde{q}}(hc_i)) = \mathcal{O}(h^{\min(2s, u)+1}).$$

Hence

$$\begin{aligned} L_{\text{exact}}(q(0), q(h); h) - L(q(0), q(h); h) &= \int_0^h \mathcal{L}(q_{\text{EL}}, \dot{q}_{\text{EL}}) dt - h \sum_i b_i \mathcal{L}(\tilde{q}(hc_i), \dot{\tilde{q}}(hc_i)) \\ &= \mathcal{O}(h^{\min(2s, u)+1}). \end{aligned}$$

By variational error analysis, this proves that the integrator is of order $\min(2s, u)$. ■

Possible extension to forced systems

External forces can be accommodated by doubling the dimension*: consider the extended Lagrangian

$$\mathcal{L}^f(q, Q, \dot{q}, \dot{Q}) = \mathcal{L}(Q, \dot{Q}) - \mathcal{L}(q, \dot{q}) + \frac{1}{2}(f(Q, \dot{Q}) + f(q, \dot{q}))(Q - q).$$

First taking variations with respect to Q , then imposing $Q = q$ we find

$$\frac{\partial \mathcal{L}(q, \dot{q})}{\partial q} - \frac{d}{dt} \frac{\partial \mathcal{L}(q, \dot{q})}{\partial \dot{q}} + f(q, \dot{q}) = 0.$$

Good news: we can use this to apply variational error analysis to forced systems.[†]

Bad news: our superconvergence result requires that critical curves of are minimizers, but this does not hold for \mathcal{L}^f .

We needed this assumption to show that the minimizing polynomial \tilde{q} is close to the polynomial interpolating the continuous solution \hat{q} .

Workaround? Without this assumption, can we still expect the difference $\hat{q} - \tilde{q}$ will be small in a generic case?

*Galley. [Classical mechanics of nonconservative systems](#). Physical review letters, 2013.

[†]De Diego, de Almagro. [Variational order for forced Lagrangian systems](#). Nonlinearity, 2018.

Conclusions

- ▶ Under mild assumptions (minimality of critical curves) Galerkin variational integrators based on polynomials of degree s are of order $2s$.
Proof is inspired on the calculus of variations itself.
- ▶ Forced systems pose an interesting challenge.
- ▶ The foundational result on variational error analysis is subtle. Can we provide a more insightful proof using modified Lagrangians?

Thank you for your attention!